On a Toy Model for Three-Dimensional Conformal Probability

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Main references

For big picture:

A.A., "Towards three-dimensional conformal probability", arXiv:1511.03180[math.PR] (27 pages).

For technical details:

A.A., A. Chandra and G. Guadagni, "Rigorous quantum field theory functional integrals over the p-adics I: anomalous dimensions", arXiv:1302.5971 [math.PR] (162 pages).

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Hence $\mathbb{T} = \bigcup_{k \in \mathbb{Z}} \mathbb{L}_k$ naturally has the structure of a doubly infinite tree which is organized into layers or generations \mathbb{L}_k :



Picture for d = 1, p = 2

Forget $[0,\infty)^d$ and \mathbb{R}^d and just keep the tree. Define the substitute for the continuum " $\mathbb{L}_{-\infty}$ " := leafs at infinity. Forget $[0,\infty)^d$ and \mathbb{R}^d and just keep the tree. Define the substitute for the continuum " $\mathbb{L}_{-\infty}$ ":= leafs at infinity.

More precisely, these are the infinite bottom-up paths in the tree.



A path representing an element $x \in \mathbb{L}_{-\infty}$

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We also define the origin $\mathbf{0} \in \mathbb{L}_0$ as the leftmost path/lattice site. Thus $\#(\bar{B}(\mathbf{0}, p^k)) = p^{dk}$ as for the volume of balls in \mathbb{R}^d .

The massless fractional Gaussian field:

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To every litter G of Mama Cat $z \in L_{k+1}$ associate a centered Gaussian random vector $(\zeta_x)_{x\in G}$ with $p^d \times p^d$ covariance matrix made of $1 - p^{-d'}$ s on the diagonal and $-p^{-d'}$ s everywhere else. We impose that Gaussian vectors corresponding to different layers or different litters are independent.

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The ancestor function: for $k \ge 0$, $\mathbf{x} \in \mathbb{L}_0$, let $\operatorname{anc}_k(\mathbf{x})$ denote the ancestor in \mathbb{L}_k . Let the symbol $[\phi]$ denote a parameter to be specified in $(0, \frac{d}{2})$.

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$$\psi_{\mathbf{x}} = \sum_{k \ge 0} p^{-k[\phi]} \zeta_{\mathrm{anc}_k(\mathbf{x})} \; .$$

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It satisfies

$$C_{\mathbf{x}\mathbf{y}} := \mathbb{E}\psi_{\mathbf{x}}\psi_{\mathbf{y}} \sim rac{\mathrm{Cst}}{d(\mathbf{x},\mathbf{y})^{2[\phi]}}$$

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at long distance. Define the probability measure μ on $\mathbb{R}^{\mathbb{L}_0}$ as the law of ψ .

$$\psi_{\mathbf{x}}^{2} := \psi_{\mathbf{x}}^{2} - C_{\mathbf{00}}$$

 and

$$:\psi_{\mathbf{x}}^{4}:=\psi_{\mathbf{x}}^{4}-6C_{\mathbf{00}}\psi_{\mathbf{x}}^{2}+3C_{\mathbf{00}}^{2}.$$

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as well as the couplings $(g, \mu) \in (0, \infty) \times \mathbb{R}$.

g is the phi-four coupling constant and μ is the mass.

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From now on: d = 3, $[\phi] = \frac{3-\epsilon}{4}$ for $\epsilon > 0$ small. We also fix $L = p^{\ell}$ for some integer $\ell \ge 1$.

For $m \geq 0$ we define the probability measure ν_m on $\mathbb{R}^{\mathbb{L}_0}$

$$d\nu_m(\phi) \sim \exp\left(-\sum_{\mathbf{x}\in \bar{B}(\mathbf{0},L^m)} \left\{g:\phi^4_{\mathbf{x}}:+\mu:\phi^2_{\mathbf{x}}:
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and the weak limit $\nu = \lim_{m \to \infty} \nu_m$.

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For g in an interval of size $\sim \epsilon^{\frac{3}{2}}$ aroung \bar{g}_* we constructed a critical mass $\mu_c(g)$ for which the infinite volume measure ν is critical, i.e., $\mathbb{E}_{\nu}\phi_{\mathbf{x}}\phi_{\mathbf{y}} \to 0$ when $d(\mathbf{x}, \mathbf{y}) \to \infty$ and $\sum_{\mathbf{x} \in \mathbb{L}_0} \mathbb{E}_{\nu}\phi_{\mathbf{0}}\phi_{\mathbf{x}} = \infty$.

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Using a new rigorous renormalization group method, we proved that in an spatially averaged sense

$$\mathbb{E}_{
u} \phi_{\mathbf{x}} \phi_{\mathbf{y}} \sim rac{ ext{Cst}}{d(\mathbf{x},\mathbf{y})^{2[\phi]}}$$

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where $N[\phi_x^2]$ denotes the recentered square field $\phi_x^2 - \mathbb{E}_{\nu}\phi_x^2$ and

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Remark: For 2D Ising, the scaling dimensions are $[\phi] = \frac{1}{8}$ (spin field) and $[\phi^2] = 1$ (energy field).

$$\sum_{\mathbf{x},\mathbf{y}\in\bar{B}(\mathbf{0},L^n)}\min\left[1,\frac{1}{d(\mathbf{x},\mathbf{y})}\right]^{\alpha}$$
$$=L^{3n}\left[1+\frac{1-p^{-3}}{1-p^{\alpha-3}}(L^{(3-\alpha)n}-1)\right]\sim\frac{1-p^{-3}}{1-p^{\alpha-3}}\times L^{(6-\alpha)n}$$

$$\sum_{\mathbf{x},\mathbf{y}\in\bar{B}(\mathbf{0},L^n)}\min\left[1,\frac{1}{d(\mathbf{x},\mathbf{y})}\right]^{\alpha}$$
$$=L^{3n}\left[1+\frac{1-p^{-3}}{1-p^{\alpha-3}}(L^{(3-\alpha)n}-1)\right]\sim\frac{1-p^{-3}}{1-p^{\alpha-3}}\times L^{(6-\alpha)n}$$

What we proved is

$$\sum_{\mathbf{x},\mathbf{y}\in\bar{B}(\mathbf{0},L^n)}\mathbb{E}_{\nu}\phi_{\mathbf{x}}\phi_{\mathbf{y}}\sim\mathrm{Cst}\times L^{(6-2[\phi])n}$$

and

$$\sum_{\mathbf{x},\mathbf{y}\in\bar{B}(\mathbf{0},L^n)} \mathbb{E}_{\nu} \mathcal{N}[\phi_{\mathbf{x}}^2] \mathcal{N}[\phi_{\mathbf{y}}^2] \sim \mathrm{Cst} \times L^{(6-2[\phi^2])n}$$

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Consistent with predictions by conformal bootsrap for 3D Ising with long-range interactions $J_{xy} = -(-\Delta_{\mathbb{Z}^3})_{xy}^{\frac{3+\epsilon}{4}}$.

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Consistent with predictions by conformal bootsrap for 3D Ising with long-range interactions $J_{xy} = -(-\Delta_{\mathbb{Z}^3})_{xy}^{\frac{3+\epsilon}{4}}$. Critical scaling limit believed to be a 3D CFT.



 \sim lower half-space model for conformal geometry.

Thank you for your attention.